

Hierarchical four-body dynamics in post-Newtonian regime

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Motion in Keplerian potential

Shape of the orbit:

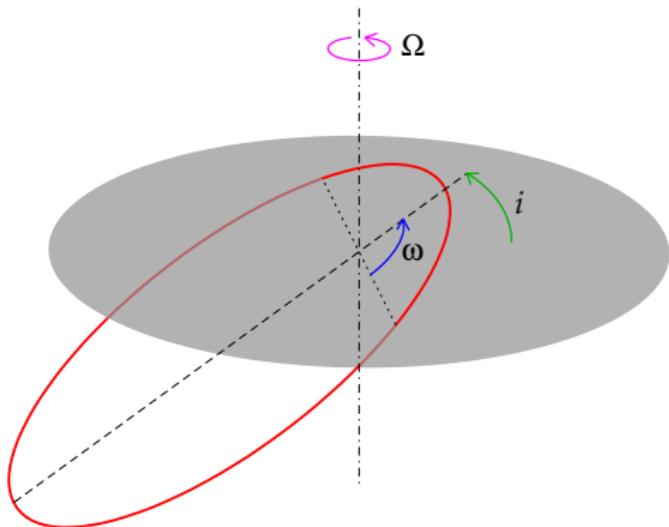
- semi-major axis, a
- eccentricity, e

Orientation of the orbital plane:

- inclination, $i \in (0^\circ, 180^\circ)$
- longitude of the ascending node, $\Omega \in (0^\circ, 360^\circ)$

Orientation and position within the orbital plane:

- argument of pericentre, ω
- true anomaly, ν

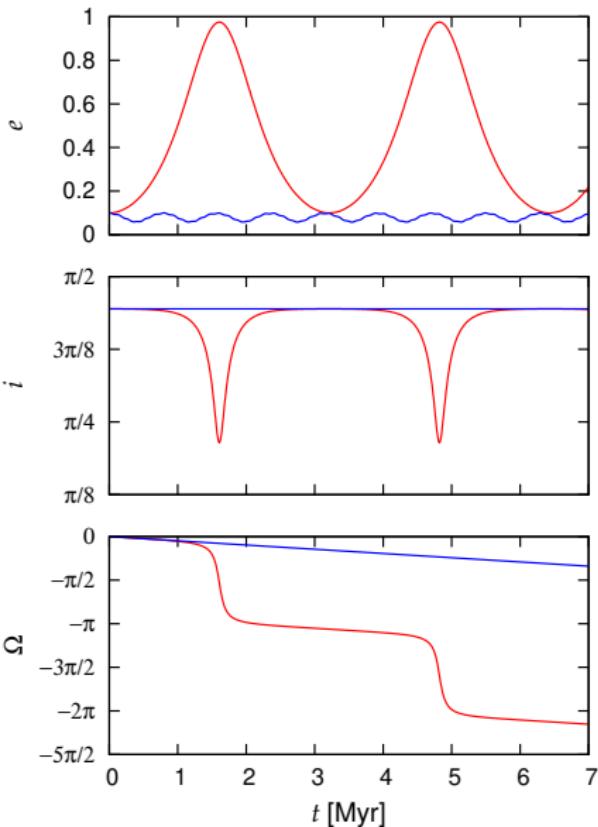


Kozai-Lidov oscillations

Broken spherical symmetry →
angular momentum is not an
integral of motion →
oscillations of eccentricity and
inclination

Example:

- $M_{\text{BH}} = 3.5 \times 10^6 M_{\odot}$
- $M_{\text{CND}} = M_{\text{BH}}$,
 $R_{\text{CND}} = 1.5 \text{ pc}$
- $a_0 = 0.1 R_{\text{CND}}$, $e_0 = 0.1$
 $i_0 = 80^\circ$, $\omega_0 = 0$, $\Omega_0 = 0$
- $M_c = 0.1 M_{\text{BH}}$



'Kozai' equations

$$T_K \sqrt{1 - e^2} \frac{de}{dt} = \frac{15}{8} e (1 - e^2) \sin 2\omega \sin^2 i$$

$$T_K \sqrt{1 - e^2} \frac{di}{dt} = -\frac{15}{8} e^2 \sin 2\omega \sin i \cos i$$

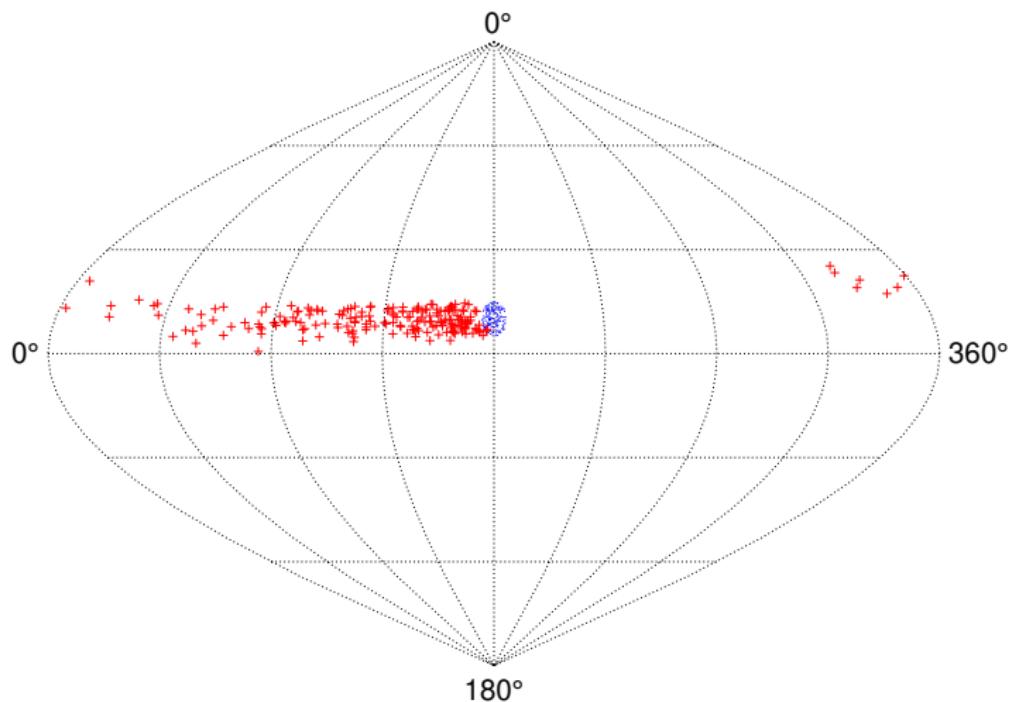
$$T_K \sqrt{1 - e^2} \frac{d\omega}{dt} = \frac{3}{4} \{2 - 2e^2 + 5 \sin^2 \omega [e^2 - \sin^2 i]\} + \left(\frac{d\omega}{dt}\right)_c$$

$$T_K \sqrt{1 - e^2} \frac{d\Omega}{dt} = -\frac{3}{4} \cos i [1 + 4e^2 - 5e^2 \cos^2 \omega]$$

$$T_K \equiv \frac{M_{\text{BH}}}{M_{\text{CND}}} \frac{R_{\text{CND}}^3}{a \sqrt{GM_{\text{BH}}a}}$$

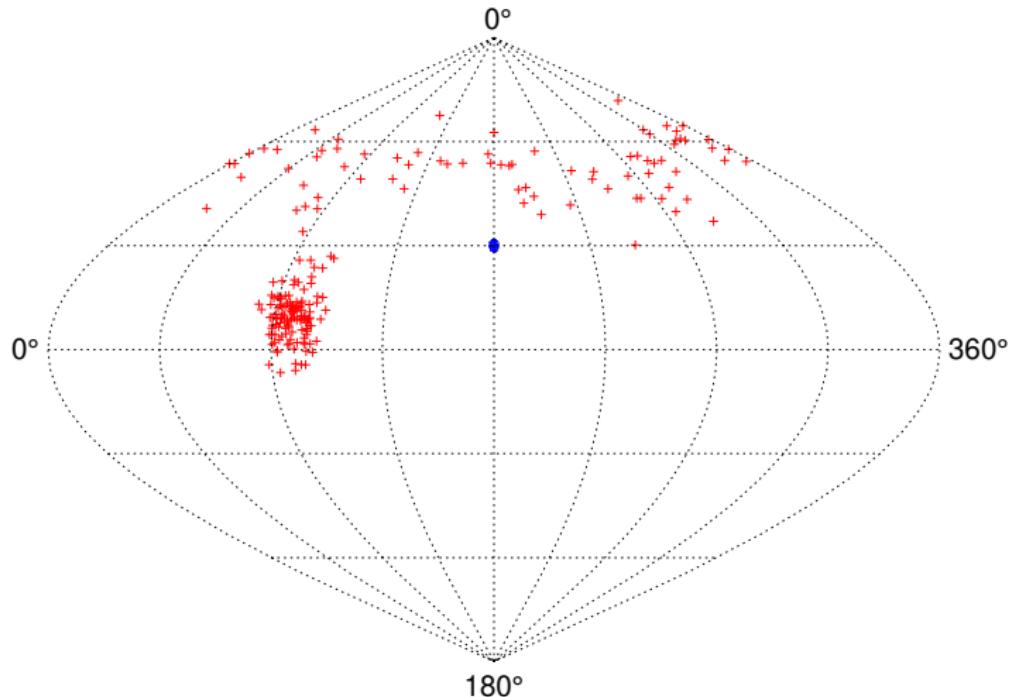
$$\frac{d\Omega}{dt} \approx -\frac{3}{4} \frac{\cos i}{T_K} \frac{1 + \frac{3}{2}e^2}{\sqrt{1 - e^2}} \approx \text{const} \text{ for } \left(\frac{d\omega}{dt}\right)_c \text{ sufficiently large}$$

Warped disc of test particles



(Šubr, Schovancová & Kroupa 2009)

Warped self-gravitating disc



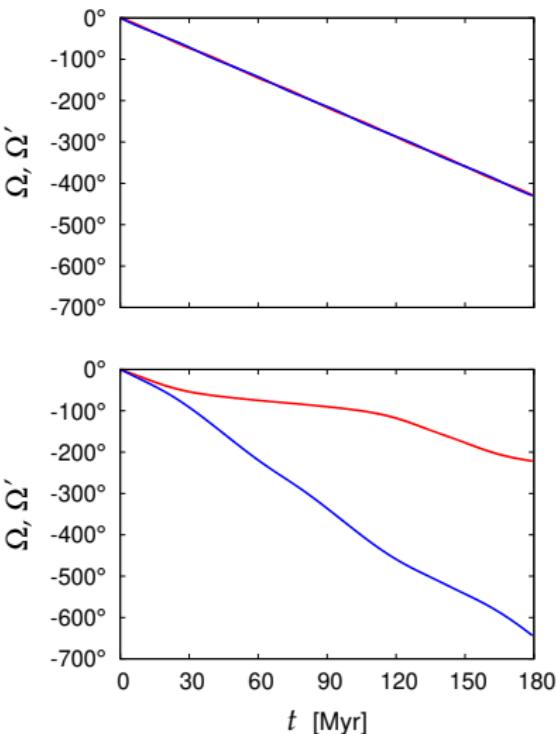
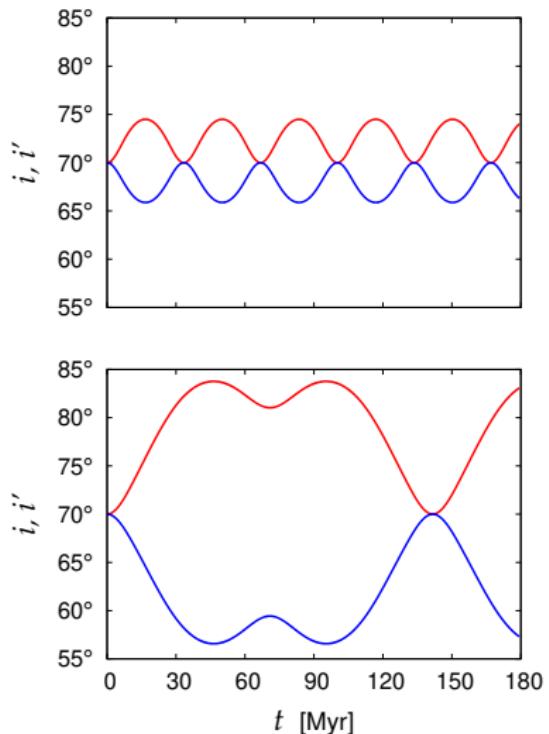
(Haas, Šubr & Kroupa 2011)

Two interacting stars in a perturbed Keplerian potential

Secular theory for evolution of two nearby orbits in a perturbed Keplerian potential developed in
Haas, Šubr & Vokrouhlický (2011).

- dominating Keplerian potential ($M_{\text{BH}} = 3.5 \times 10^6 M_{\odot}$)
- perturbing body on circular orbit ($M_{\text{CND}} = 0.3 M_{\text{BH}}$,
 $R_{\text{CND}} = 1.5 \text{ pc}$)
- extended spherical potential damping the Kozai-Lidov oscillations
- two light bodies on initially nearby orbits
 - $i = i' = 70^\circ$
 - $\Omega = \Omega' = 0$
 - $e = e' = 0$
 - $a = 0.05 R_{\text{CND}}$, $a' = 0.04 R_{\text{CND}}$
 - $m = m' = 9 \times 10^{-6} M_{\text{BH}} \vee 5 \times 10^{-6} M_{\text{BH}}$

Two interacting stars in a perturbed Keplerian potential

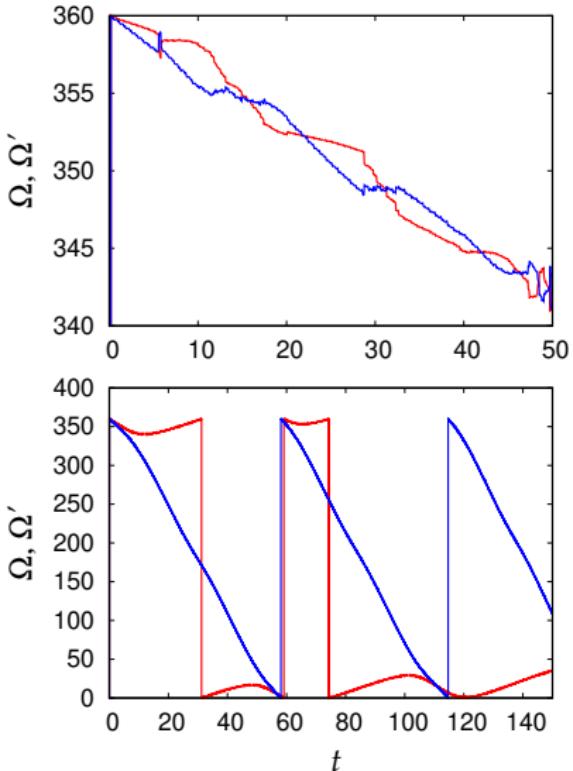
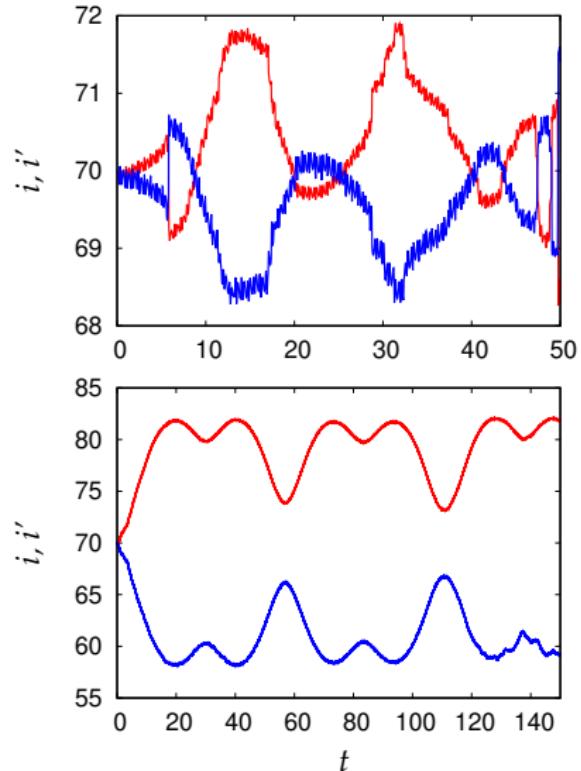


(Haas, Šubr & Vokrouhlický 2011)

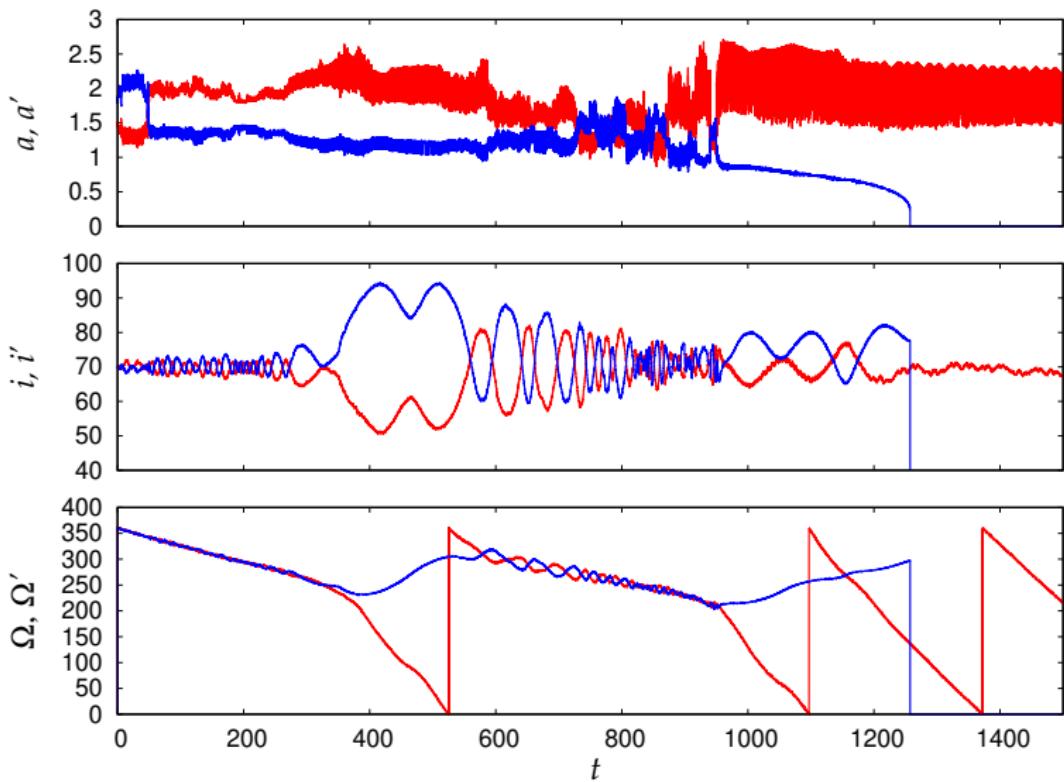
Two stars in perturbed post-Newtonian potential

Direct integrations of equations of motion by means of ARCHAIN code (Chassonney, Capuzzo-Dolcetta & Mikkola, 2019; algorithm by Mikkola & Merritt, 2006, 2008), involving post-Newtonian terms up to the 2.5 order.

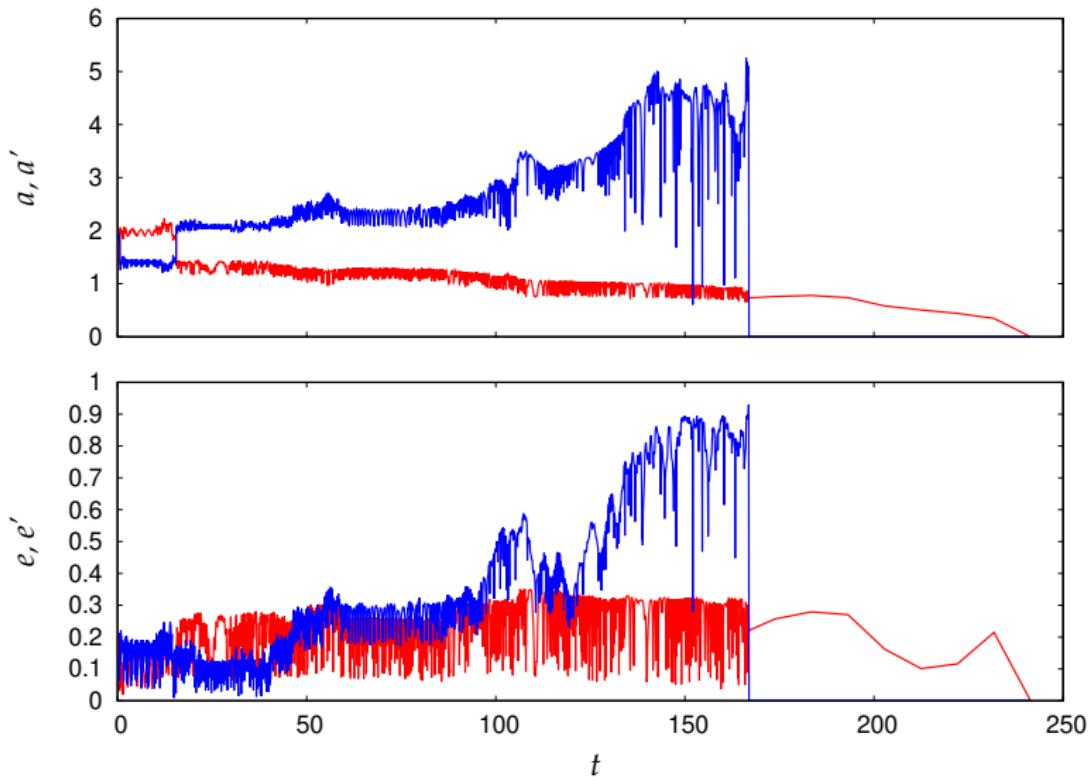
Two stars in perturbed post-Newtonian potential



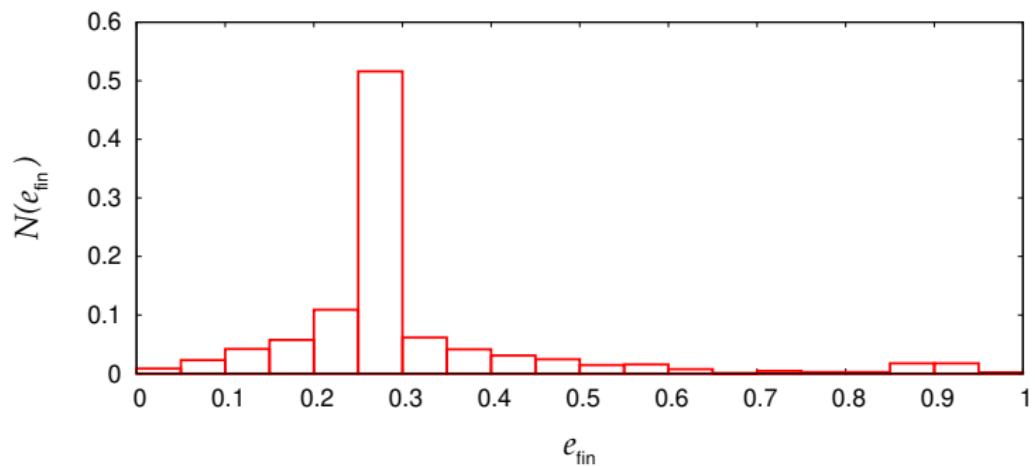
Swap, re-binding, inspiral



High eccentricity inspiral?



High eccentricity inspiral?



Summary

- hierarchical four-body dynamics in P-N exhibits similar modes of evolution to the Newtonian case (with arbitrary damper of Kozai-Lidov oscillations)
- mutual interaction of ‘light’ bodies often accelerates inspiral
 - energy transfer?
 - eccentricity oscillations?
 - effectively larger mass?
- inspiral with large (> 0.8) eccentricity in non-negligible number ($\approx 4\%$) of cases
- results very preliminary with arbitrary parameters of the system
 - revision needed
 - seek for astrophysically relevant application

High eccentricity inspiral?

